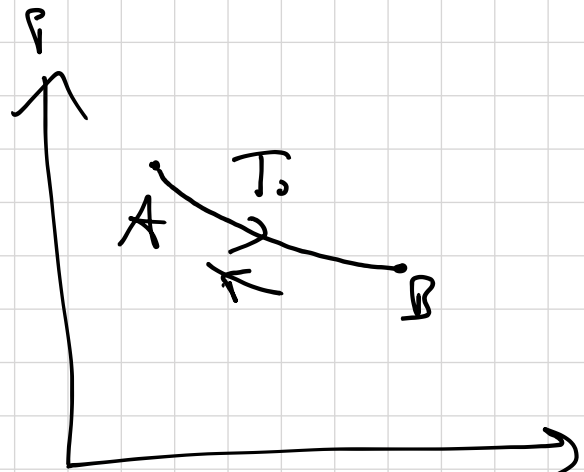
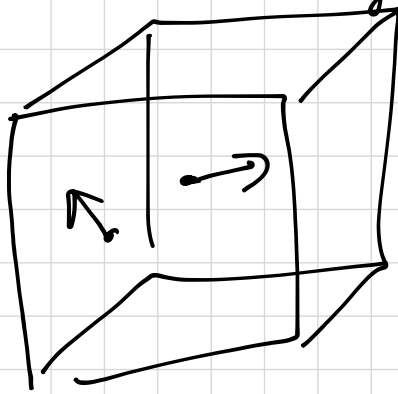


reversible change



isothermal expansion of an ideal gas: V

$$W_{A \rightarrow B} = \int_A^B P dV = \int_{V_A}^{V_B} \frac{nRT_0}{V} dV =$$

$$nRT_0 \log \frac{V_B}{V_A} > 0 \text{ if } V_B > V_A$$

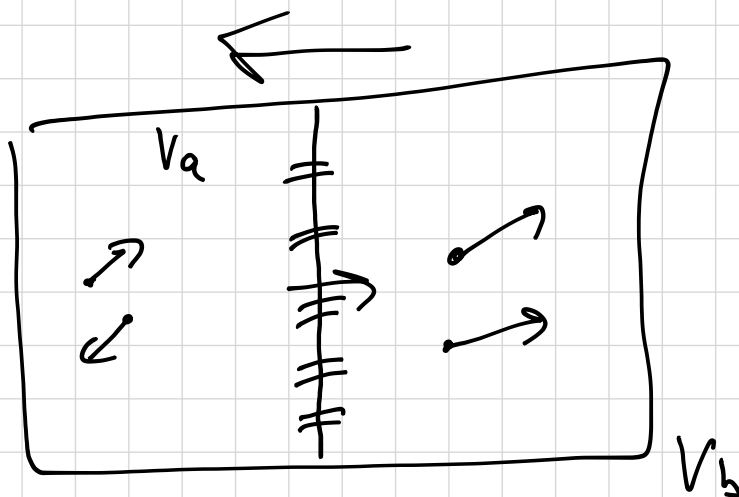


Because T_0 is constant, $\Delta E(T_0) = 0$,

$$\Delta E = Q - W$$

$$Q_{A \rightarrow B} = W_{A \rightarrow B}$$

irreversible change - Joule (free) expansion



$$\begin{aligned} Q &= 0 \\ W &= 0 \end{aligned}$$

$$\Delta E = 0 //$$

$$\Delta T = 0 //$$

Reversible and irreversible processes

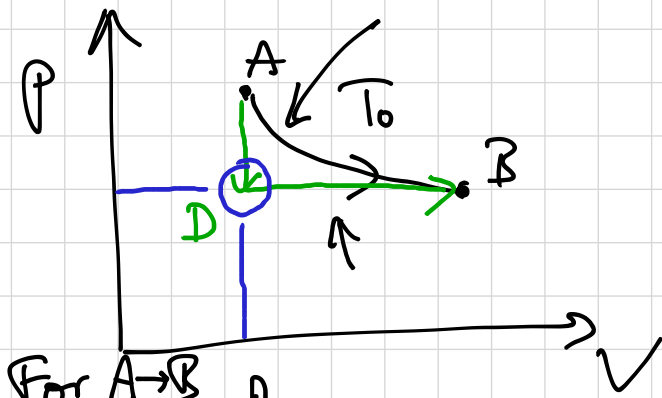
Entropy

$$\Delta S = \frac{Q}{T} \quad \text{if } T \text{ does not change in the process}$$

more generally

$$\Delta S = \int_A^B \frac{dQ}{T}$$

Entropy is a state variable!



For $A \rightarrow B$

$$\Delta S = \int_A^B \frac{dQ}{T_0} = \frac{Q_{A \rightarrow B}}{T_0} = \frac{nRT_0 \log \frac{V_B}{V_A}}{T_0}$$

$$= nR \log \frac{V_B}{V_A} \rightarrow 0$$

For $A \rightarrow D$

$$\Delta V = 0$$

$$Q = nC_v \Delta T$$

$$\Delta S = \int_A^D \frac{dQ}{T} = \int_A^D \frac{nC_v dT}{T} = nC_v \int_A^D \frac{dT}{T}$$

$$= nC_v \log \frac{T_D}{T_A} \leftarrow$$

For $D \rightarrow B$

$$\Delta S = \int_D^B \frac{dQ}{T} = \int_D^B \frac{nC_p dT}{T} = nC_p \log \frac{T_B}{T_D}$$

$$= n(C_v + R) \log \frac{T_B}{T_D} \leftarrow$$

For $A \rightarrow D \rightarrow B$:

$$\Delta S = nC_v \log \frac{T_D}{T_0} + n(C_v + R) \log \frac{T_0}{T_D}$$

$$= \cancel{nC_v} \log \frac{T_D}{T_0} + \cancel{nC_v} \log \frac{T_0}{T_D} + nR \log \frac{T_0}{T_D}$$

$$= nR \log \frac{T_0}{T_D}$$

$$pV = nRT$$

$$nRT_0 = p_A V_A = p_B V_B$$

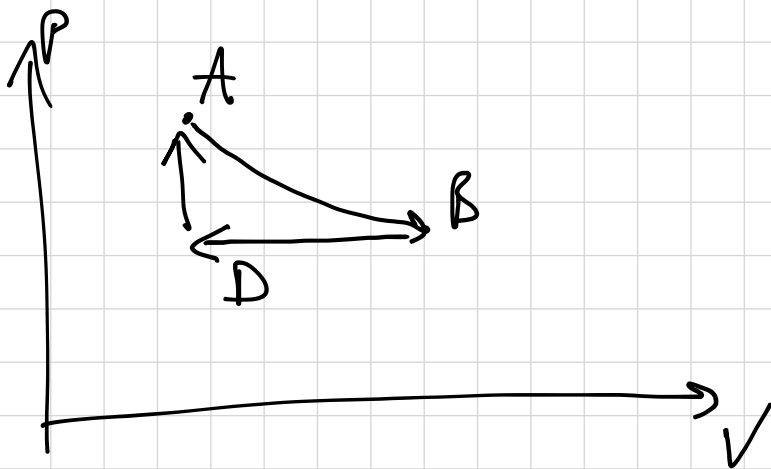
$$nRT_D = p_B V_A$$

$$= nR \log \frac{P_B V_B \cancel{nR}}{P_B V_A \cancel{nR}}$$

$$\Delta S_{A \rightarrow D \rightarrow B} = nR \log \frac{V_B}{V_A} > 0$$

$$\underline{\Delta S_{A \rightarrow B}} = \underline{\Delta S_{A \rightarrow D \rightarrow B}} \quad \checkmark$$

→ S (entropy) is a state variable



$$S_{A \rightarrow B} + S_{B \rightarrow D} + S_{D \rightarrow A} = 0$$

- $S_{A \rightarrow D \rightarrow B}$

$$\oint \frac{dQ}{T} = 0$$

Comparing the isothermal and the free expansion,

$$\left. \begin{aligned} \Delta S_{\text{reversible}} &= nR \log \frac{V_B}{V_A} > 0 \\ \Delta S_{\text{irreversible}} &= 0 \end{aligned} \right\}$$

And in general,

$$\Delta S = \underbrace{\int_A^B \frac{dQ_{\text{reversible}}}{T}}_{>} \underbrace{\int_A^B \frac{dQ_{\text{irreversible}}}{T}}$$

Example:

w/o
friction

w/
friction

Second law of Thermodynamics
thermally isolated systems:

$$\Delta S \geq 0$$

Third law of Thermodynamics

$$T \rightarrow 0 \quad \underline{S \rightarrow 0} \leftarrow$$

$$T \rightarrow 0$$

$$\underline{S \rightarrow S_0} \leftarrow$$

Entropy and the first laws of thermodynamics.

for a reversible process, we can write

$$dS = \frac{dQ}{T}, \quad dQ = TdS \quad dW = PdV$$

$$dE = TdS - PdV \leftarrow$$

$$dE = \underbrace{dQ}^{\#} - \underbrace{dW}^{\#} \leftarrow$$

Statistical interpretation of entropy

$$dE = TdS - PdV,$$

$$\begin{aligned} & \uparrow \\ & = \frac{\partial E}{\partial S} \Big|_V \end{aligned} \quad \frac{1}{T} = \frac{\partial S}{\partial E} \Big|_V \quad *$$

$$\frac{d \ln \Omega(E)}{dE} = \frac{1}{k_B T} \quad *$$

← # of microstates

$$k_B \frac{d \ln \Omega}{dE} = \frac{\partial S}{\partial E} \Big|_V$$

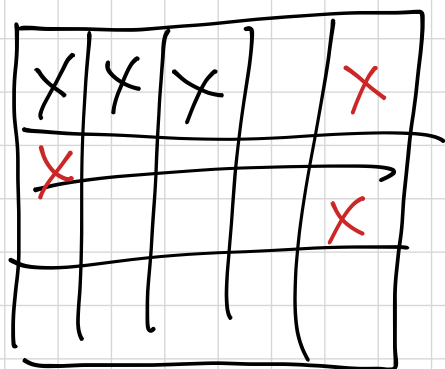
$$S = k \ln \Omega \leftarrow$$

Since the 2nd law of thermodynamics

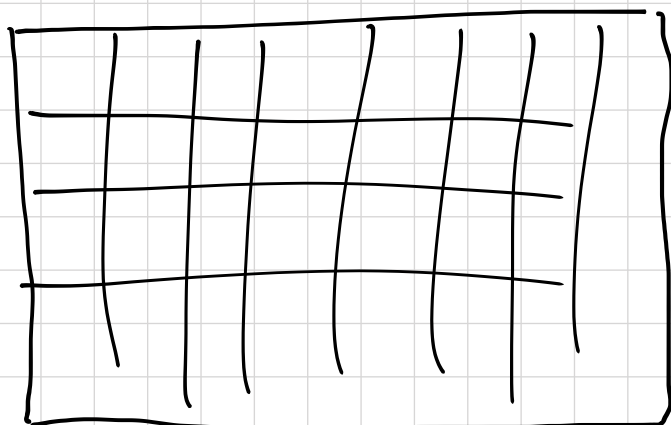
$$\Delta S \geq 0,$$

more order \rightarrow less order
(fewer microstates) (more microstates)

Free expansion:



m spots



$$S = k_B \ln \Omega$$

$$\Omega = \frac{m!}{(m-N)! N!} = \binom{m}{N}$$

sterling's formula: $\ln x! = x \ln x - x$

$$S = k_B \ln \frac{m!}{(m-N)! N!}$$

$$= k_B \left[\ln m! - \ln (m-N)! - \ln N! \right]$$

$$\approx k_B \left[m \ln m - \cancel{m} - (m-N) \ln (m-N) + \cancel{m-N} - N \ln N + \cancel{N} \right]$$

$$= k_B \left[m \ln m - (m-N) \ln (m-N) - N \ln N \right]$$

$$= k_B \left[\underbrace{m \ln m} - \underbrace{m \ln (m-N)} + \underbrace{N \ln (m-N)} - \underbrace{N \ln N} \right]$$

$$= k_B \left[-m \ln \frac{m-N}{m} + N \ln \frac{m-N}{N} \right]$$

$$= k_B \left[-m \ln \left(1 - \frac{N}{m} \right) + N \ln \left(\frac{m}{N} - 1 \right) \right]$$

$\downarrow \quad \quad \quad \downarrow \quad \uparrow$
 $0 \quad \quad \quad \infty$

m spär, $m \rightarrow \infty$

$$\approx k_B \left[-m \left(-\frac{N}{m} \right) + N \ln \frac{m}{N} \right]$$

$$= k_B \left(N + N \ln \frac{m}{N} \right)$$

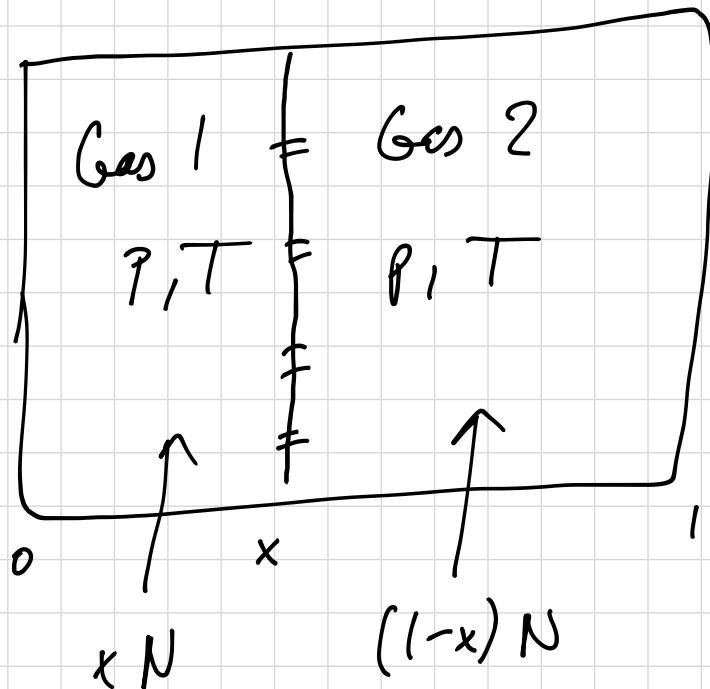
$\Delta S : m \rightarrow (2m)$

$$\begin{aligned}
 S(2m) - S(m) &= k_B \left(\cancel{N} + N \ln \frac{2m}{N} \right) \\
 &\quad - k_B \left(\cancel{N} + N \ln \frac{m}{N} \right) \\
 &= k_B N \ln \left(\frac{2m}{N} \frac{N}{m} \right)
 \end{aligned}$$

$$= k_B N \ln 2$$

$$\Delta S = k_B N \log \frac{V_b}{V_a} \Big|_{V_b=2V_a} = k_B N \log 2$$

Entropy of mixing



$$\Delta S = \Delta S_1 + \Delta S_2 = k_B x N \log \left(\frac{1}{x} \right) + k_B (1-x) N \log \left(\frac{1}{1-x} \right)$$

$$= k_B N \left[-x \log x - (1-x) \log (1-x) \right]$$

